

W 35-196



Hosenbüchel W 35-196

Gastherrie

IV Hans Fleming

$u = u_1 + u_2 + \dots + u_n$   
 $v = v_1 + v_2 + \dots + v_n$   
 $w = w_1 + w_2 + \dots + w_n$

$$F(a_1, \dots, a_n) = \int_{x_1(a_1, \dots, a_n)}^{x_2(a_1, \dots, a_n)} \dots \int_{x_n(a_1, \dots, a_n)} dx_1 \dots dx_n$$

$$= \int_{\alpha_1}^{\alpha_2} \dots \int_{\alpha_n} \frac{\partial(x_1, \dots, x_n)}{\partial(\alpha_1, \dots, \alpha_n)} d\alpha_1 \dots d\alpha_n$$

1. Phasenschicht - Bewegungsgleichungen  
 $\rho(p, q, t)$  - wogende Fläche  $e, c, \dots$  - Ray  
 präzisionspunkt  $\alpha_0, \beta_0, \gamma_0, \dots$   
 $\alpha = \alpha_0 + \dots$   
 $\beta = \beta_0 + \dots$   
 $\gamma = \gamma_0 + \dots$   
 die selben Relationen  $S, p, q, c$   
 $\alpha, \beta, \gamma = 1, p, c$   
 Die Substitution der Elemente durch  
 die Elementare ist linear mit koeff.  $p, q$  abhängig  
 daher kann die Determinante aus dem  $J$  abgelesen  
 haben werden

$$\frac{dQ}{dt} = \frac{\partial Q}{\partial t} - \left[ \frac{\partial p u Q}{\partial x} + \frac{\partial p v Q}{\partial y} + \dots \right] + \rho \left[ x \frac{\partial Q}{\partial u} + y \frac{\partial Q}{\partial v} + \dots \right]$$

$$\frac{\partial Q}{\partial t} = - \left( \frac{\partial p u}{\partial x} + \dots \right) \quad Q = 1 \text{ Kontinuitätsgleichung}$$

$$u = \bar{u} + \bar{u}' \quad v = \bar{v} + \bar{v}' \quad w = \bar{w} + \bar{w}'$$

$\bar{u}, \bar{v}, \bar{w}$  betrachtet  $\bar{u}', \bar{v}', \bar{w}'$  molekular

$$\rho \frac{\partial Q}{\partial t} - Q \left( \frac{\partial p u}{\partial x} + \frac{\partial p v}{\partial y} + \dots \right) = \rho \frac{\partial Q}{\partial t} - \left[ \frac{\partial p u Q}{\partial x} + \frac{\partial p v Q}{\partial y} + \dots \right] + \rho \left[ x \frac{\partial Q}{\partial u} + y \frac{\partial Q}{\partial v} + \dots \right]$$

$$\frac{\partial Q}{\partial u} = \frac{\partial Q}{\partial \bar{u}} \quad \frac{\partial Q}{\partial v} = \frac{\partial Q}{\partial \bar{v}}$$

$$\frac{\partial (p u Q)}{\partial x} = p u \frac{\partial Q}{\partial x} + Q \frac{\partial (p u)}{\partial x}$$

$$\rho \left[ \frac{\partial \bar{u}}{\partial t} + \bar{u} \frac{\partial \bar{u}}{\partial x} + \bar{v} \frac{\partial \bar{u}}{\partial y} + \bar{w} \frac{\partial \bar{u}}{\partial z} \right] = \rho \frac{D \bar{u}}{D t} - \left[ \frac{\partial}{\partial x} \rho \bar{u} \bar{u} + \frac{\partial}{\partial y} \rho \bar{u} \bar{v} \right] + \rho \left[ \bar{u} \frac{\partial \bar{u}}{\partial x} + \dots \right]$$

$\rho \bar{u} = \rho u$   
 in 2 Zonen  $\rho$   $u_1, u_2, \dots, u'_1, u'_2$

$$\bar{u} u = \bar{u}(\bar{u} + \bar{u}') = \bar{u}^2 + \bar{u} \bar{u}' \quad u_1 + u_2 = u'_1 + u'_2 \quad \text{er} \quad \frac{D \bar{u}}{D t} = 0$$

$$\rho \frac{d \bar{u}}{d t} = - \frac{\partial (\rho \bar{u}^2)}{\partial x} - \frac{\partial (\rho \bar{u} \bar{v})}{\partial y} - \frac{\partial (\rho \bar{u} \bar{w})}{\partial z} + \rho X$$

$$\rho \frac{d \bar{v}}{d t} = - \frac{\partial (\rho \bar{u} \bar{v})}{\partial x} - \frac{\partial (\rho \bar{v}^2)}{\partial y} - \frac{\partial (\rho \bar{v} \bar{w})}{\partial z} + \rho Y$$

$$\rho \frac{d \bar{w}}{d t} = - \frac{\partial (\rho \bar{u} \bar{w})}{\partial x} - \frac{\partial (\rho \bar{v} \bar{w})}{\partial y} - \frac{\partial (\rho \bar{w}^2)}{\partial z} + \rho Z$$

$f_x = \rho X$   
 $e / \text{Kor}$   
 $u \text{ u } \bar{u}$

$$\rho \frac{d u}{d t} = \rho X - \frac{\partial X_x}{\partial x} - \frac{\partial X_y}{\partial y} - \frac{\partial X_z}{\partial z} \quad \text{no } \frac{d u}{d t} \text{ b } \frac{D u}{D t}$$

$X_x = \rho \bar{u}^2 \quad Y_y = \rho \bar{v}^2 \quad Z_z = \rho \bar{w}^2$   
 wenn man diese Annahmen macht  $\rho \bar{u} = \rho u = \text{erl } \bar{u}$   
 $X_y = Y_x = \rho \bar{u} \bar{v} \quad Y_z = Z_y = \rho \bar{v} \bar{w} \quad Z_x = X_z = \rho \bar{u} \bar{w}$

$\rho \bar{u} \bar{v} = \rho \bar{v} \bar{u}$   
 in Hydrodynamik  $e \quad X_x = Y_y = Z_z = \rho$   
 $X_x = \dots Z_z = 0$

nicht reichendes  $\rho$   $\bar{u} \bar{v} = \bar{v} \bar{u} = \bar{u} \bar{v} = 0$

$$\rho \bar{u}^2 = \rho \bar{v}^2 = \rho \bar{w}^2 = \rho = \frac{1}{3} \rho \bar{c}^2$$

$$\rho = \frac{\rho \bar{u} \bar{c}^2}{3} \quad \text{Gasdruckformel}$$

$$\rho \frac{D \bar{u}}{D t} = 0 - \rho \bar{u}^2 + \bar{u}^2 + \bar{v}^2 + \bar{w}^2$$

in Hydrodynamik  $X_y = \left( \frac{\partial \rho u}{\partial y} - \frac{\partial \rho v}{\partial x} \right) e \text{ u } \bar{u}$   
 $\bar{u} \bar{v} = \bar{v} \bar{u}$

$\rho \bar{u} \bar{v} = \rho \bar{v} \bar{u}$   
 $\rho \bar{u} \bar{w} = \rho \bar{w} \bar{u}$

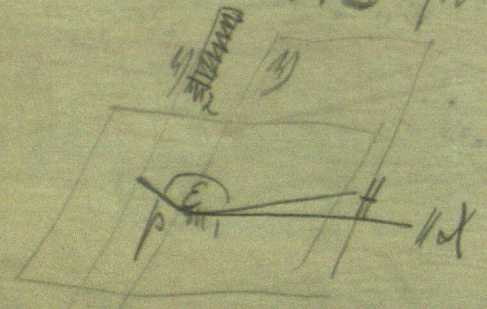
$\varphi(x) = 0 \quad x > 0$   
 $\varphi(x) > 1 \quad x \leq 0$

Comp:  $0 < \varphi(x) < 1$

Zusammenstoß  $p$  &  $w$  in  $x$  & marktliche Wechselwirkung stattfindet.

1.  $dw_1$
2.  $dw_2$

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Klasse  $k_1$   $dw_1$   
 $k_2$   $dw_2$  } Comp  $p \& w$

Gattung I:  
 $k_1 \& k_2$   $\epsilon + \delta \epsilon$   
 $p + \delta p$

$f \in p \& w$

Comp  $p$  &  $w$  1 & 2 primativen  $\& p \& w$  de  $\& w$ .

$p \& w \in \mathcal{M} f_1 dw_1 f_2 dw_2 \in p \& w$

$Q_1 = Q(u, v, w_1) \text{ ig } \mathcal{M} p$

$\mathcal{M} \text{ Co } Q_1 + Q_2 \quad \mathcal{M} \text{ Co } Q'_1 + Q'_2$

$\& w$   $\& w$   $p \& w$   $\& w$   $[Q'_1 + Q'_2 - Q_1 + Q_2]$

$\& w$   $p \& w$   $\& w$   $p, \epsilon u v w u, v, w_1$

$\frac{dQ}{dt} \cdot dt = \frac{1}{2} \int_0^{\infty} \int_0^{\infty} p \& w \& w [Q'_1 + Q'_2 - Q_1 + Q_2]$

$\epsilon \& w$   $\& w$

Klassen  $k'_1$   $dw'_1$   $f'_1$   
 $k'_2$   $dw'_2$   $f'_2$

$\left. \begin{matrix} Q_1 \\ Q_2 \end{matrix} \right\} \rightarrow \begin{cases} \epsilon_1 \\ \epsilon_2 \end{cases}$

$\left. \begin{matrix} Q'_1 \\ Q'_2 \end{matrix} \right\} \rightarrow \begin{cases} \epsilon_1 \\ \epsilon_2 \end{cases}$

$$N \frac{dQ}{dt} = \frac{1}{2} \left[ \iint d\omega_1 \iint d\omega_2 \left( \rho' d\epsilon \right) \left( \frac{1}{r_{12}} \right) (Q_1 + Q_2 - (Q_1 + Q_2)) \right]$$

$\downarrow$   $\mathcal{H}I$

o o t was  $m_2$  z e Partelsymmetrie:



$\uparrow$   $\mathcal{H}I$

u r o z z f r d k r e n t e u p e s m o r e

u  $d\omega_1, d\omega_2$  r z u d p r o f t o r u d  $d\omega, d\omega_2 = d\omega_1, d\omega_2$   
 u f u o f o r o r n e E n e r g i e p r i n z, s ~ F l ä c h e n m a ß

$$S \frac{dQ}{dt} = \frac{1}{4} \left[ \iint d\omega_1 \iint d\omega_2 \left( \rho \rho' d\epsilon \right) \left( \frac{1}{r_{12}} \right) (Q_1 + Q_2) \left( \frac{1}{r_{12}} \right) \right]$$

$\rho \omega \omega \omega / \omega_1 \omega_2 \omega_3 \omega_4 \omega_5 \omega_6 \omega_7 \omega_8 \omega_9 \omega_{10}$   
 $\varphi(x) / 12 \text{ z o r n n } \} \text{ m o f o r m u l e}$

$$m_1 \frac{d^2 x_1}{dt^2} = \varphi(x) \frac{x_1 - x_2}{r} \quad m_2 \frac{d^2 x_2}{dt^2} = \varphi(x) \frac{x_2 - x_1}{r}$$

o p r e i v u g y g r o

$$\frac{d^2(x_2 - x_1)}{dt^2} = \left( \frac{1}{m_1} + \frac{1}{m_2} \right) \varphi(x) \frac{x_2 - x_1}{r}$$

$$x_2 - x_1 = \xi$$

$$\frac{1}{M}$$

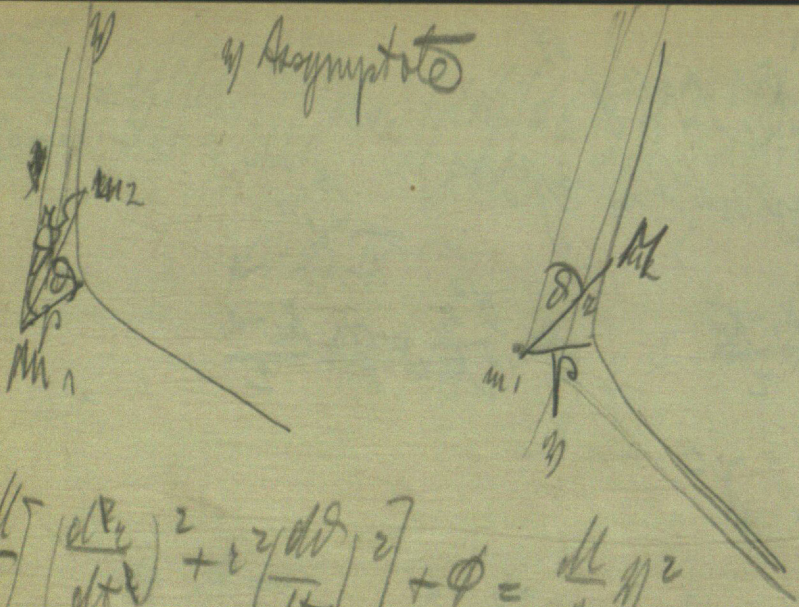
Zweikörperproblem.

$$M \frac{d^2 \xi}{dt^2} = \varphi(x) \frac{\xi}{r} \quad M \frac{d^2 \eta}{dt^2} = \varphi(x) \frac{\eta}{r}$$

e n p r e i v u g y g r o 20 M.

~~u r o z z f r d k r e n t e u p e s m o r e~~  
 e r u g u n z e p o r t r e z f r o

4) Asymptote



$$\frac{dL}{dt} \left[ \left( \frac{dr}{dt} \right)^2 + r^2 \left( \frac{d\vartheta}{dt} \right)^2 \right] + \Phi = \frac{dL}{dt} \eta^2$$

$$\frac{1}{2} r^2 \frac{d\vartheta}{dt} = \text{const}$$

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$$\varphi = - \int_{\infty}^r \varphi(r) dr$$

$$\varphi = - \frac{\partial \Psi}{\partial r}$$

$\Psi$  pot. w. eff.

$$\frac{dL}{dt} \left[ \left( \frac{dr}{dt} \right)^2 + r^2 \left( \frac{d\vartheta}{dt} \right)^2 \right] + \Psi = \frac{dL}{dt} \eta^2 \quad \Psi(\infty) = 0$$

$$\frac{1}{2} r^2 \frac{d\vartheta}{dt} = \frac{1}{2} \eta^2 p$$

$$dt = \frac{r^2 d\vartheta}{\eta p}$$

$$\frac{dL}{dt} \left[ \frac{dr^2 + r^2 d\vartheta^2}{r^4 d\vartheta^2} \cdot \eta^2 p^2 \right] + \Psi = \frac{dL}{dt} \eta^2$$

$$p^2 \left[ \frac{1}{r^4} \left( \frac{dr}{d\vartheta} \right)^2 + \frac{1}{r^2} \right] + \frac{2\Psi}{dL \eta^2} = 1$$

$$\frac{p}{r} = g \quad \vartheta$$

$$\frac{p^2}{r^4} dr^2 = \left( +1 - \frac{2\Psi}{dL \eta^2} - \frac{p^2}{r^2} \right) d\vartheta^2$$

$$- \frac{p}{dr} = dg$$

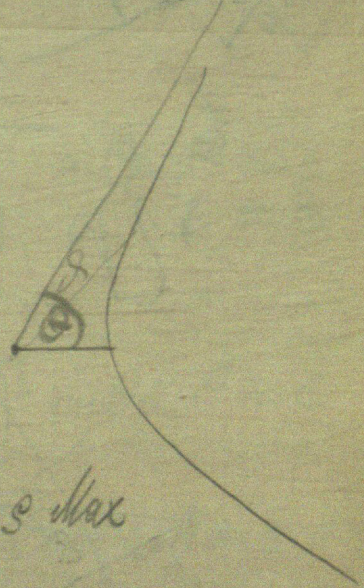
$$dg^2 = \left( 1 - g^2 - \frac{2\Psi}{dL \eta^2} \right) d\vartheta^2$$

$$g = \int_{\vartheta_0}^{\vartheta} \frac{1}{\sqrt{1 - g^2 - \frac{2\Psi}{dL \eta^2}}} d\vartheta$$

$$\vartheta = \int_0^{g_{\text{max}}} \frac{dg}{\sqrt{1 - g^2 - \frac{2\Psi}{dL \eta^2}}}$$

$$\sqrt{1 - g^2 - \frac{2\Psi}{dL \eta^2}} = 0 \quad g_{\text{max}}$$

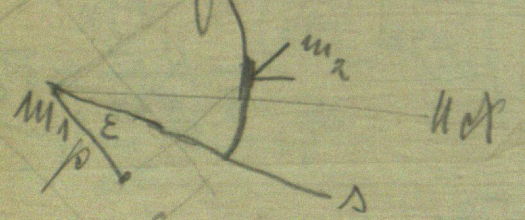
Spezialwertung  $\varphi = \frac{K}{r^{n+1}}$



$$\varphi = \frac{k}{mv^2} = \frac{hc}{mv^2}$$

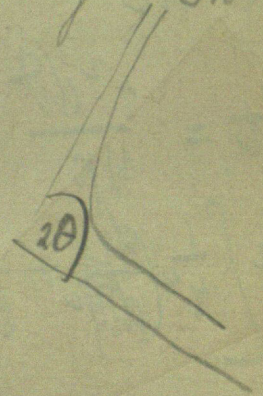
$$\Theta = \int_0^{\theta_{max}} \frac{dq}{\sqrt{1 - \beta^2 - \frac{2kq}{mv^2}}}$$

Abstand  $r(\pi - \Theta)$

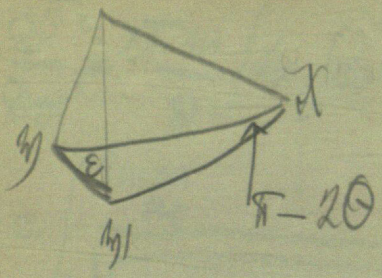
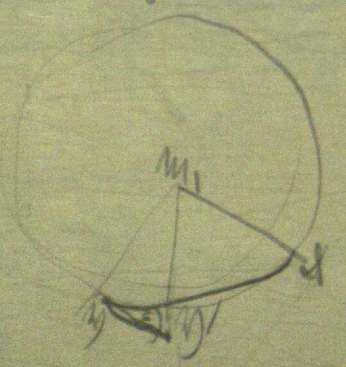


$$\epsilon = \pi - \left\{ \rho, (m) \right\}$$

$e \sin \Theta \approx \sin 2\Theta$



$\sin \theta \approx \theta$  in  $m \bar{x} \bar{y}$



$$\cos(\eta'x) = \cos(\eta x) \cos(\pi - 2\Theta) + \sin(\eta x) \sin(\pi - 2\Theta) \cos \epsilon$$

$$\begin{aligned} & \eta' \cos(\eta'x) \text{ relative to } \eta \text{ in } \bar{L}_0 \\ & = u_2' - u_1' = (u_2 - u_1) \cos 2\Theta + \sqrt{v^2 - (u_2 - u_1)^2} \sin 2\Theta \cos \epsilon \end{aligned}$$

$$-\cos 2\Theta = 1 - 2\cos^2 \Theta$$

$$u_2' - u_1' = (u_2 - u_1) - 2(u_2 - u_1) \cos^2 \Theta + 2 \sqrt{v^2 - (u_2 - u_1)^2} \sin \Theta \cos \Theta \cos \epsilon$$

$$u_2' + u_1' = u_2 + u_1$$

$$u_2' = u_2 - (u_2 - u_1) \cos^2 \Theta + \sqrt{v^2 - (u_2 - u_1)^2} \sin \Theta \cos \Theta \cos \epsilon$$

$$u_1' = u_1 + (u_2 - u_1) \cos^2 \Theta - \sqrt{v^2 - (u_2 - u_1)^2} \sin \Theta \cos \Theta \cos \epsilon$$

$$e = \frac{d\Theta}{d\epsilon} \quad w = u_2 + \frac{d\Theta}{d\epsilon}$$

$\int \dots \int \dots$   
 $Q = n^2$

$$\underbrace{n_1^2 - n_1^2 + n_2^2 - n_2^2}$$

$$\begin{aligned}
 &= (n_2 - n_1) \cos^2 \theta + n^2 (n_2 - n_1) \sin^2 \theta \cos \epsilon \\
 &+ 2 n_1 (n_2 - n_1) \cos^2 \theta + (\dots) \cos \epsilon \\
 & \text{1. \& 2. \& 3. \& 4.}
 \end{aligned}$$

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$$\frac{dQ}{dt} = \frac{m}{2} \int \dots \int \dots (Q_1 + Q_2 - Q_1 - Q_2)$$

$$\frac{dQ}{dt}$$

$$\begin{aligned}
 n_1^2 + n_2^2 - n_1^2 - n_2^2 &= 2 \sin^2 \theta \cos^2 \theta [(n_2 - n_1)^2 + (\dots)] \\
 \int d\epsilon(\theta) &= 2 \sin^2 \theta \cos^2 \theta \pi [n^2 - 3(n_2 - n_1)^2] + (\dots) \cos \epsilon
 \end{aligned}$$

$$\frac{dQ}{dt} = \frac{\pi m}{2} \int \dots \int \dots [n^2 - 3(n_1 - n_2)^2] \dots$$

$$\theta = \int_0^{\theta_{\max}} \frac{d\theta}{\sqrt{1 - s^2 - \frac{2k}{m} \frac{g^n}{u}}}$$

$$s_{\max} / p^n = \sqrt{u} = 0$$

$$\xi^n = p^n \frac{m n^2}{k}$$

$$\theta = \int_0^{\theta} \frac{d\theta}{\sqrt{1 - s^2 - \frac{2k}{m} \frac{g^n}{u}}}$$

$$p d p = g d g \left( \frac{k}{m n^2} \right)^{\frac{1}{n}}$$

$$\frac{dQ}{dt} = \pi m \int \dots \int \dots [n^2 - 3(n_1 - n_2)^2] \left( \frac{k}{m n^2} \right)^{\frac{1}{n}} \dots$$

weiter Spezialisierung  $n=4$

$$C \cdot e \left( \frac{1}{n^2} \right)^{\frac{1}{n}} \text{ sehr unlegitim?}$$

$$\text{keine Zahl} = \frac{13682}{48} = \frac{A}{48}$$

$$a \cdot C \cdot \int \dots \int \dots \sqrt{1 - s^2} \dots$$



$$u^2 \rho \frac{d\bar{n}^2}{dt} = \frac{\rho u^2}{4} \sqrt{\frac{2k}{M}} A [\sin(\omega_1 t + \phi_1) + \sin(\omega_2 t + \phi_2)]^2 - 3(\bar{n}_1 - \bar{n}_2)^2$$

$$y^2 = (\bar{n}_1 - \bar{n}_2)^2 + (m_1 - m_2)^2 + (m_1 - m_2)^2$$

$$\int \sin \omega_1 t \int \sin \omega_2 t \bar{n}_1^2 = N \int \sin \omega_1 t \bar{n}_1^2 = N^2 \bar{n}^2$$

$$\int \sin \omega_1 t \int \sin \omega_2 t \bar{n}_1 \bar{n}_2 = N \bar{n} \bar{n} = 0 \quad M = \frac{m}{2} \rho$$

$$\rho \frac{d\bar{n}^2}{dt} = \frac{m}{4} \sqrt{\frac{2k}{M}} A [\bar{v}^2 - \bar{n}^2] \quad \bar{v}^2 = \bar{n}^2 + \bar{m}^2 + \bar{m}^2$$

$$u = \bar{u} + \bar{n} \quad \text{I}$$

$$\rho \frac{d\bar{u}}{dt} = \rho \lambda - \frac{\partial}{\partial x} (\rho \bar{n}^2) - \frac{\partial}{\partial y} (\rho \bar{n} \bar{m}) - \frac{\partial}{\partial z} (\rho \bar{n} \bar{m})$$

nicht reichende Messen:  $\bar{n} \bar{m} = \bar{n} \bar{m} = \bar{m} \bar{m} = 0$   
 $\bar{m}^2 = \bar{m}^2 = \bar{m}^2 = \frac{c^2}{3} = \frac{p}{\rho}$

wir wollen die Gleichungen, reichende Messen:

$$\rho \frac{d\bar{Q}}{dt} = \rho \frac{d\bar{Q}}{dt} + \frac{\partial}{\partial x} (\rho u \bar{Q}) + \frac{\partial}{\partial y} (\rho u \bar{Q}) + \dots - \rho \lambda \frac{\partial \bar{Q}}{\partial u} + \dots$$

f3 = 7415 exakt

$$c = u \rho \frac{d\bar{Q}}{dt} \quad \bar{Q} = u^2 \bar{n} \bar{n} \dots \frac{c^2}{3} \rho$$

$$\bar{Q} = u^2 = \bar{u}^2 + 2\bar{u}\bar{n} + \bar{n}^2 \quad \bar{Q} = \bar{u}^2 + \bar{n}^2$$

$$\bar{n} \bar{Q} = \bar{n} \bar{u}^2 + 2\bar{u} \bar{n}^2 + \bar{n}^3 = 2\bar{u} \bar{n}^2$$

$$\bar{v} \bar{Q} = 0 \quad \bar{w} \bar{Q} = 0 \quad \frac{\partial \bar{Q}}{\partial \bar{u}} = 2\bar{u}$$

$$\rho \frac{d}{dt} \bar{n}^2 = \rho \frac{d}{dt} (\bar{u}^2 + \bar{n}^2) + \frac{\partial}{\partial x} (2\bar{u} \rho \bar{n}^2) - \rho \lambda 2\bar{u}$$

a 2 p p^2, x / y / z = ...

$$\rho \frac{d\bar{u}}{dt} = 2p \frac{\partial \bar{u}}{\partial x} + \rho \frac{d}{dt} p$$

$$\rho \frac{d}{dt} (\bar{n}^2 + \bar{m}^2 + \bar{m}^2) = 2p \left( \frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{u}}{\partial y} + \frac{\partial \bar{u}}{\partial z} \right) + \rho \frac{d}{dt} p$$

= 0 p u / el p Co / 1 r / c

# Statistische Mechanik

von Maxwell, Boltzmann, Gibbs

Die kinetische Theorie eine quadratische Form der Geschwindigkeiten

$$d \ln \Omega = \rho \frac{d \bar{p}}{p} = \rho \bar{v}$$

$$\rho \frac{d \bar{u}}{dt} = \frac{2}{3} \rho \left[ 2 \frac{\partial \bar{u}}{\partial x} - \frac{\partial \bar{v}}{\partial y} - \frac{\partial \bar{w}}{\partial z} \right] + [\epsilon]$$

Correspondence  $\rho \frac{d \bar{u}}{dt}$  f. ...

$$[\epsilon^2 - 3 \bar{u}^2] = \frac{4}{m} \sqrt{\frac{m}{2k}} \frac{1}{A} \frac{2}{3} \rho \left[ 2 \frac{\partial \bar{u}}{\partial x} - \frac{\partial \bar{v}}{\partial y} - \frac{\partial \bar{w}}{\partial z} \right] + [\epsilon]$$

...  $\rho \frac{d \bar{u}}{dt} = \dots$

...  $\rho \frac{d \bar{u}}{dt} = \dots$

$$\rho \bar{u}^2 = \rho \frac{2p}{\Delta \rho} \sqrt{\frac{2m^3}{k}} \left\{ 2 \frac{\partial \bar{u}}{\partial x} - \frac{\partial \bar{v}}{\partial y} - \frac{\partial \bar{w}}{\partial z} \right\}$$

$\frac{2R}{2}$

$$\rho \frac{du}{dt} = \rho X - \frac{\partial p}{\partial x} - R \left[ \Delta \bar{u} + \frac{1}{3} \frac{\partial}{\partial x} \left( \frac{\partial \bar{u}}{\partial x} + \frac{\partial \bar{v}}{\partial y} + \frac{\partial \bar{w}}{\partial z} \right) \right]$$

$$\mathcal{L}(q, \dot{q}) = \bar{\mathcal{L}}(q, p)$$

$$\bar{\mathcal{L}} = \frac{1}{2m} p_1^2 + \frac{1}{m\omega^2} p_2^2$$

$$\frac{\partial \mathcal{L}}{\partial q_i} \Big|_{q, p} = - \frac{\partial \mathcal{L}}{\partial q_i} \Big|_{q, \dot{q}}$$

$$\frac{\partial \bar{\mathcal{L}}}{\partial p_i} \Big|_p = - \frac{1}{m\omega^2} p_2^2$$

$$\frac{\partial \mathcal{L}}{\partial p_i} \Big|_{q, p} = \dot{q}_i$$

$$\frac{\partial \bar{\mathcal{L}}}{\partial q} = 0$$

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial t} \Big|_p &= - \frac{1}{m\omega^2} m\omega^2 \dot{q}^2 \\ &= -m\omega^2 \dot{q}^2 = - \frac{\partial \mathcal{L}}{\partial t} \Big|_{q, \dot{q}} \end{aligned}$$

$$\mathcal{L} = \frac{m}{2} (\dot{r}^2 + r^2 \dot{\varphi}^2)$$

$$q_1 = r \quad \frac{\partial \mathcal{L}}{\partial r} = m r \dot{\varphi}^2$$

$$q_2 = \varphi \quad \frac{\partial \mathcal{L}}{\partial \varphi} = 0$$

$$p_1 = \frac{\partial \mathcal{L}}{\partial \dot{r}} = m \dot{r}$$

$$p_2 = \frac{\partial \mathcal{L}}{\partial \dot{\varphi}} = m r^2 \dot{\varphi}$$

$$\dot{r} = \frac{1}{m} p_1 \quad \dot{\varphi} = \frac{1}{m r^2} p_2$$

$$\bar{\mathcal{L}} = \frac{m}{2} \left( \frac{1}{m^2} p_1^2 + \frac{r^2}{m^2 r^4} p_2^2 \right)$$

$$\frac{\partial \bar{\mathcal{L}}}{\partial p_1} \Big|_{p, q} = \frac{1}{m} p_1 = \dot{r}$$

$$\frac{\partial \bar{\mathcal{L}}}{\partial p_2} = \frac{1}{m r^2} p_2 = \dot{\varphi}$$

$f = \text{rot} \circ p \circ \text{rot} = h$   
 Hamiltonsche & f.z.h.  
 $\langle \text{rot} \circ p \circ \text{rot} \rangle$  / Hamiltonsche  
 v.

$$Q_i = - \frac{\partial \mathcal{V}}{\partial q_i} \quad \mathcal{V} = \mathcal{V}(q_i, t)$$

$q_i$  - allgem. Koordinate  
 $p_i$  - Impuls

$\bar{V}$  Lage Konfiguration  
 65.0.0.0.0.0.0.0

$$-\frac{\partial \bar{V}}{\partial \bar{q}_i} = + \frac{d p_i}{dt} - \frac{\partial \bar{L}}{\partial \bar{q}_i}$$

$$= + \frac{d p_i}{dt} + \frac{\partial \bar{L}}{\partial \bar{q}_i}$$

$$\frac{d p_i}{dt} = - \frac{\partial (\bar{V} + \bar{L})}{\partial \bar{q}_i}$$

$$\frac{d \bar{q}_i}{dt} = \frac{\partial \bar{L}}{\partial p_i} = \frac{\partial (\bar{L} + \bar{V})}{\partial p_i}$$

in  $\frac{\partial \bar{V}}{\partial p_i} = 0$   $\bar{V} + \bar{L} = E$

$$\frac{d p_i}{dt} = - \frac{\partial E}{\partial \bar{q}_i} \quad \frac{d \bar{q}_i}{dt} = \frac{\partial E}{\partial p_i}$$

0 0 0  $E = E(q_i, p_i)$   
 Hamilton's  $\bar{H} = \bar{L}$

$\partial \bar{L} / \partial p_i = \dot{\bar{q}}_i$

$e \frac{\partial \bar{V}}{\partial \bar{q}_i} = 0$  in separable  
 in  $\bar{V}$

$E(q, p) = \text{const } e_i$

$$\frac{dE}{dt} = \sum \left[ \frac{\partial E}{\partial p_i} \frac{d p_i}{dt} + \frac{\partial E}{\partial q_i} \frac{d \bar{q}_i}{dt} \right] = 0$$

für  $e_i$

as as:

$\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

0 in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

in  $\bar{V}$   $\bar{L}$   $\bar{H}$   $\bar{E}$   $\bar{L} + \bar{V}$

$$\Delta = \frac{\partial(p_1, \dots, p_n, q_1, \dots, q_n)}{\partial(p'_1, \dots, p'_n, q'_1, \dots, q'_n)}$$

0.1.1910.

$$E(q, p, t) = \bar{L} + \bar{V}$$

$$q_i = q_i(t, q_i^0, \dots, q_n^0, p_i^0, \dots, p_n^0)$$

$$\frac{\partial(q_1, \dots, q_n, p_1, \dots, p_n)}{\partial(q_1^0, \dots, q_n^0, p_1^0, \dots, p_n^0)} = \frac{\partial(q, p)}{\partial(q^0, p^0)}$$

$$\Delta = \frac{\partial(p', q')}{\partial(p, q)}$$

$$q_i' = q_i + \frac{\partial q_i}{\partial t} dt = q_i + \dot{q}_i dt$$

$$= q_i + \frac{\partial E}{\partial p_i} dt$$

$$p_i' = p_i + \dot{p}_i dt = p_i - \frac{\partial E}{\partial q_i} dt$$

$$q_1' = q_1 + \frac{\partial E}{\partial p_1} dt$$

$$\frac{\partial q_1'}{\partial q_1} = 1 + \frac{\partial^2 E}{\partial q_1 \partial p_1} dt \quad \frac{\partial q_1'}{\partial q_2} = \frac{\partial^2 E}{\partial q_2 \partial p_1} dt$$

$$\frac{\partial q_1'}{\partial p_1} = \frac{\partial^2 E}{\partial p_1^2} dt, \quad \frac{\partial q_1'}{\partial p_n} = \frac{\partial^2 E}{\partial p_n \partial p_1} dt$$

$$\Delta = \begin{vmatrix} 1 + \frac{\partial^2 E}{\partial p_1 \partial p_1} dt & \frac{\partial^2 E}{\partial p_1 \partial p_2} dt & \dots & \frac{\partial^2 E}{\partial p_1 \partial p_n} dt \\ \frac{\partial^2 E}{\partial p_2 \partial p_1} dt & 1 + \frac{\partial^2 E}{\partial p_2 \partial p_2} dt & \dots & \frac{\partial^2 E}{\partial p_2 \partial p_n} dt \\ \dots & \dots & \dots & \dots \\ \frac{\partial^2 E}{\partial p_n \partial p_1} dt & \dots & \dots & 1 + \frac{\partial^2 E}{\partial p_n \partial p_n} dt \end{vmatrix}$$

$$q_2' = q_2 + \frac{\partial E}{\partial p_2} dt$$

$$p_1' = p_1 - \frac{\partial E}{\partial q_1} dt$$

can be written as  $\sum \frac{\partial^2 E}{\partial p_i \partial p_j} dt$

$$1 + \frac{\partial^2 E}{\partial p_1 \partial p_1} dt + \frac{\partial^2 E}{\partial p_1 \partial p_2} dt + \dots + \frac{\partial^2 E}{\partial p_1 \partial p_n} dt$$

$$- \frac{\partial^2 E}{\partial p_1 \partial q_1} dt - \dots - \frac{\partial^2 E}{\partial p_1 \partial p_n} dt$$

$$\Delta = 1 \quad \frac{\partial(q' p')}{\partial(q, p)} = 1 \quad \text{and} \quad \frac{\partial(q'' p'')}{\partial(q', p')} = 1$$

$$\frac{\partial(q'' p'')}{\partial(q, p)} = \frac{\partial(q'' p'')}{\partial(q', p')} \cdot \frac{\partial(q', p')}{\partial(q, p)} = 1 \cdot 1 = 1$$

and  $p_1, p_2, \dots, p_n, q_1, q_2, \dots, q_n = 1$

$$\text{we } \frac{\partial(q, p)}{\partial(q, p)} = 1$$

can be written as the same system of equations.  $E(p, q)$  is the same as  $E(p_0, q_0)$  and  $\int \dots \int dp_1 \dots dp_n dq_1 \dots dq_n = G_0$

$$L = \int \dots \int dp_{11} dq_{11}$$

$$\int \dots \int dp_{11} dq_{11} = \int \dots \int dp_{11}^0 dq_{11}^0 \Delta(p, q)$$

$$= \int \dots \int dp_{11}^0 dq_{11}^0 \quad \underline{g = g_0 \text{ Linuorles } g}$$

z. B.  $\dots$

$$L = \frac{m}{2} \dot{x}^2 \quad v = -mgx$$

$$E = \frac{1}{2m} p^2 - mgq$$

$$\frac{dp}{dt} = -\frac{\partial E}{\partial q} = mg$$

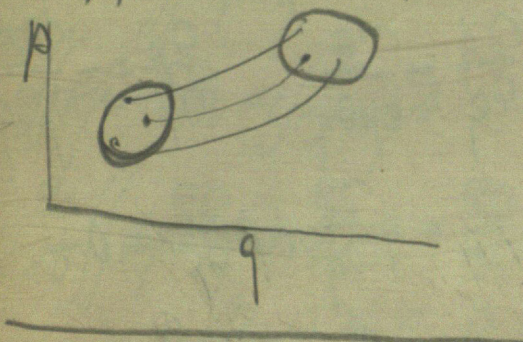
$$\frac{\partial q}{\partial t} = \frac{\partial E}{\partial p} = \frac{p}{m}$$

$$q = x \quad p = \frac{\partial L}{\partial \dot{x}} = m\dot{x}$$

$$p = p_0 + mg t$$

$$q = q_0 + \int \frac{p}{m} dt = q_0 + \frac{p_0}{m} t + \frac{g}{2} t^2$$

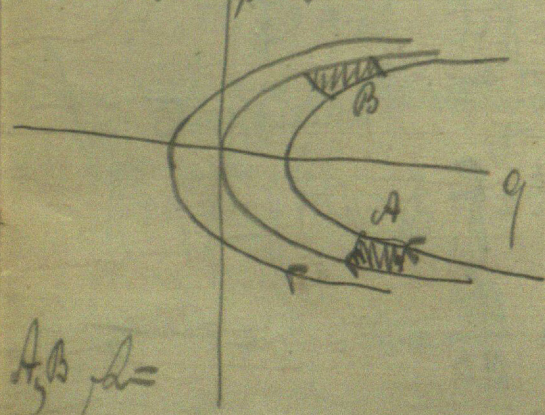
$$\frac{\partial |q, p|}{\partial |q_0, p_0|} = \begin{vmatrix} 1 & \frac{p_0}{m} \\ 0 & 1 \end{vmatrix} = 1$$



g. V. 1910.

$$E = \frac{1}{2m} p^2 - mgq$$

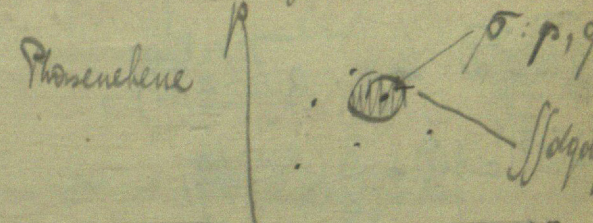
$$q = \frac{1}{2m^2 g} p^2 - \frac{E}{mg}$$



A, B,  $\dots$

$$L = \frac{m}{2} \dot{x}^2 \quad v = \dot{x}$$

$$q = x \quad p = m\dot{x}$$

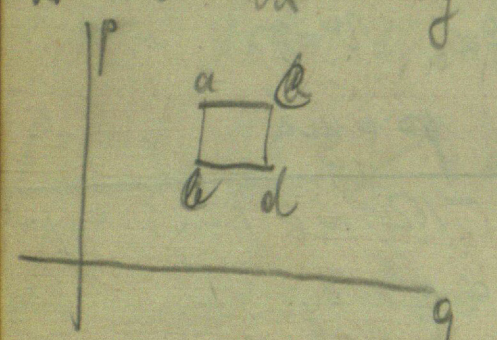


Phasenebene  $\dots$

$\dots$

$$\frac{\partial q}{\partial t} = \frac{\partial E}{\partial p} = \frac{p}{m}$$

$$\frac{d}{dt} \left( \frac{\partial p}{\partial \dot{q}} + \frac{\partial(uq)}{\partial \dot{x}} + \frac{\partial(vq)}{\partial \dot{y}} \right) = 0$$



$$\overline{ab} \dot{q} \dot{p} = \overline{ab} (\dot{q} \dot{p})'$$

$$(\dot{q} \dot{p})' = \dot{q} \dot{p} + \frac{\partial}{\partial t} (\dot{q} \dot{p}) dt$$

$$ab [\dot{q} \dot{p} - (\dot{q} \dot{p})'] = - \overline{ab}$$

$$= - dq dp \frac{\partial}{\partial q} (\dot{q} \dot{p})$$

$$- dq dp \frac{\partial}{\partial p} (\dot{q} \dot{p})$$

$$\frac{\partial \dot{p}}{\partial t} + \frac{\partial(\dot{q} \dot{p})}{\partial q} + \frac{\partial(\dot{q} \dot{p})}{\partial p} = 0$$

$\dot{p} \in \mathbb{R}$   $\dot{q} \in \mathbb{R}$

$\dot{p} \in \mathbb{R}$   $\dot{q} \in \mathbb{R}$

$$\frac{dq}{dt} = \frac{\partial E}{\partial p} \quad \frac{dp}{dt} = - \frac{\partial E}{\partial q}$$

$$\frac{\partial \dot{q}}{\partial q} = \frac{\partial^2 E}{\partial p \partial q} \quad \frac{\partial \dot{p}}{\partial p} = - \frac{\partial^2 E}{\partial p \partial q}$$

$$\frac{\partial \dot{p}}{\partial t} + \dot{q} \frac{\partial \dot{p}}{\partial q} + \dot{p} \frac{\partial \dot{p}}{\partial p} = 0$$

$$\dot{p} = \dot{p}(q, p, t) = \dot{p}$$

$$\frac{d\dot{p}}{dt} = 0 \quad (\text{Wasser-Öl})$$

das ist die Erhaltung  $\epsilon$

$$f_1(p, q, t) = \epsilon_1$$

$$f_2(\quad) = \epsilon_2$$

$$\dot{p} = f(f_1(q, p, t), f_2(q, p, t))$$

$$\frac{d\dot{p}}{dt} = 0 \quad \text{die Liouville's}$$

$$D_0 = \int \int p \dot{q} \dot{p} dq dp$$

$$D_{\dots} = \int \int p \dot{q} \dot{p} dq dp$$

$$\int \int p \dot{q} \dot{p} dq dp D_0 = \int \int p \dot{q} \dot{p} dq dp D$$

$$\frac{d}{dt} \int \int p \dot{q} \dot{p} dq dp = 0$$

$$\frac{d}{dt} (D \dot{q}) = 0$$

$$= \frac{dD}{dt} \dot{q} + \frac{d\dot{q}}{dt} D = 0$$

Liouville

11. V 1910

Liouville:  $f$  die Erhaltung der Phasenausdehnung.

$$\frac{dD}{dt} = 0$$

$$\frac{D}{N} = P = e^{\gamma} \frac{1}{g}$$

$$\frac{\partial(p \dot{q} \dot{p})}{N} dq dp = \dots$$

$\gamma$  age ist

$\gamma$  age ist

$$\epsilon \cdot \frac{\partial \dot{p}}{\partial t} = 0$$

systematisches = 10

$$E(q, p) = \text{const}$$

$$D = F(E)$$

$$-\frac{\partial \dot{p}}{\partial t} = \frac{dE}{dt} \left[ \dot{q} \frac{\partial E}{\partial q} + \dot{p} \frac{\partial E}{\partial p} \right]$$

23. V 1910.

$$H = \int \dots \int D dq_1 \dots dp_n$$

$$\frac{\partial D}{\partial t} = 0 \quad \text{w} \quad D = D(E)$$

$\int \dots \int P_i$   
 $P = e^{\frac{\varphi-E}{\theta}}$  hamonische

Phasevariable  
 $\int \dots \int e^{\frac{\varphi-E}{\theta}} dq_1 \dots dp_n = 1$

$\circ$   $\mathcal{H} = \mathcal{H}(x, p, t)$

$$\int \dots \int e^{\frac{\varphi-E}{\theta}} dq_1 \dots dp_n$$

$\circ$   $\mathcal{H} = \mathcal{H}(x, p, t)$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m})$

$$\int \dots \int e^{\frac{\varphi_2 - E_2}{\theta}} dq_{n+1} \dots dp_{n+m}$$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$$W = w_1 \cdot w_2 = \int \dots \int dq_1 \dots dp_n e^{\dots}$$

$$\int \dots \int dq_{n+1} \dots dp_{n+m} e^{\dots} = \int \dots \int dq_1 \dots dp_n e^{\dots}$$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$$\bar{q} = \int \dots \int q e^{\frac{\varphi-E}{\theta}} dq_1 \dots dp_n$$

$$b_1, \dots, b_n \quad m_1, m_2, \dots, m_n$$

$$x_1, y_1, z_1, \dots, x_n, y_n, z_n$$

$$L = \sum_i \frac{m_i}{2} (\dot{x}_i^2 + \dots)$$

$$p_i = m_i \dot{x}_i$$

$$\int \dots \int e^{\frac{\varphi-E}{\theta}} dx_1 dy_1 \dots dz_n$$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$$E = \varphi + \sum \frac{m_i}{2} (\dot{x}_i^2 + \dots)$$

$$e^{\frac{\varphi-E}{\theta}} = e^{\frac{\varphi-b}{\theta}} \cdot e^{-\frac{m_i}{2} \dot{x}_i^2}$$

$$P_i = M \int \dots \int e^{\frac{\varphi-b}{\theta}} dx_1 \dots dz_n \int e^{-\frac{m_i}{2} \dot{x}_i^2} dx_i \int e^{-\frac{m_i}{2} \dot{y}_i^2} dy_i \dots$$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$$L = \frac{m_i}{2} \dot{x}_i^2 + \dots$$

$$\frac{m_i}{2} \dot{x}_i^2 \quad L = \mathcal{H}(x, p, t)$$

$$= M \int \dots \int e^{\frac{\varphi-b}{\theta}} dx_1 \dots dz_n$$

$$\int_{-\infty}^{+\infty} \frac{m_i}{2} \dot{x}_i^2 e^{-\frac{m_i}{2} \dot{x}_i^2} \int \dots \int e^{\dots} dx_1 \dots dz_n$$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$\mathcal{H} = \mathcal{H}(q_1, \dots, q_{n+m}, p_1, \dots, p_{n+m})$

$$P_i = M \int \dots \int e^{\frac{\varphi-b}{\theta}} dx_1 \dots dz_n \int e^{-\frac{m_i}{2} \dot{x}_i^2} dx_i \int e^{-\frac{m_i}{2} \dot{y}_i^2} dy_i \dots$$



$$M \int \dots \int e^{-\frac{m \sum v_i^2}{2\theta}} dx_1 \dots dx_n = 1$$

and the same

$$\frac{m_1 \overline{x_1^2}}{2} = \frac{m_1}{2} \frac{\int_{-\infty}^{\infty} x_1^2 e^{-\frac{m_1 x_1^2}{2\theta}} dx_1}{\int_{-\infty}^{\infty} e^{-\frac{m_1 x_1^2}{2\theta}} dx_1}$$

$$\int_{-\infty}^{\infty} e^{-2x^2} dx = \sqrt{\frac{\pi}{2}} \quad \int_{-\infty}^{\infty} x^2 e^{-2x^2} dx = \frac{1}{2} \sqrt{\frac{\pi}{2}}$$

$$\frac{m_1 \overline{x_1^2}}{2} = \frac{m_1}{2} \frac{\frac{1}{2} \sqrt{\frac{\pi}{2}} \frac{2\theta}{m_1}}{\frac{1}{2} \sqrt{\frac{\pi}{2}} \frac{2\theta}{m_1}} = \frac{m_1}{4} \frac{2\theta}{m_1} = \frac{\theta}{2}$$

$e^{\frac{1}{2} \theta} e^{-\frac{m \sum v_i^2}{2\theta}} = e^{-\frac{3}{2} \theta}$   
 $\frac{3}{2} \theta = \frac{m}{2} \theta = L \quad \underline{\underline{L = \frac{m}{2} \theta}}$

w<sup>t</sup> d p q u m g r k w l d  
 24/2 1910.  
 w<sup>t</sup> d p q u m g r k w l d

w<sup>t</sup> d p q u m g r k w l d  
 $\frac{1}{2} \theta = \left( \sqrt{\frac{2\theta}{m_1}} \right)^2$

$(1/\theta)^{\frac{3V}{2}} (m_1 \dots m_n)^{\frac{3}{2}} \int \dots \int e^{-\frac{m \sum v_i^2}{2\theta}} dx_1 \dots dx_n$

$L = U + L$   
 $\int \dots \int e^{-\frac{m \sum v_i^2}{2\theta}} dq_1 \dots dq_n$   
 $= \int \dots \int e^{-\frac{m \sum v_i^2}{2\theta}} dq_1 \dots dq_n \int \dots \int e^{-\frac{m \sum v_i^2}{2\theta}} dp_1 \dots dp_n$

$L = A_1 p_1^2 + A_2 p_2^2 + \dots$   
 $\frac{\partial L}{\partial q_i} = \dots$

$L = \frac{1}{2} k_1 x_1^2 + \frac{1}{2} k_2 x_2^2 + \dots$

$\frac{\partial p_1}{\partial x_1} = \dots$   
 $\frac{\partial p_2}{\partial x_2} = \dots$

$\frac{\partial p_1}{\partial x_1} = \frac{\partial L}{\partial x_1}$   
 $\frac{\partial p_2}{\partial x_2} = \frac{\partial L}{\partial x_2}$

$\frac{\partial p_1}{\partial x_1} = \frac{\partial L}{\partial x_1}$   
 $\frac{\partial p_2}{\partial x_2} = \frac{\partial L}{\partial x_2}$

$\frac{\partial p_1}{\partial x_1} = \frac{\partial L}{\partial x_1}$   
 $\frac{\partial p_2}{\partial x_2} = \frac{\partial L}{\partial x_2}$

$\frac{\partial p_1}{\partial x_1} = \frac{\partial L}{\partial x_1}$   
 $\frac{\partial p_2}{\partial x_2} = \frac{\partial L}{\partial x_2}$

$$\int \dots e^{L^0} \dots$$

$$= \frac{D(p_1, \dots, p_n)}{D(q_1, \dots, q_n)}$$

$$L^0 = L^0(q) \quad e^{L^0} =$$

$$= \begin{vmatrix} \frac{\partial^2 L}{\partial q_1^2} & \frac{\partial^2 L}{\partial q_1 \partial q_2} & \dots \\ \frac{\partial^2 L}{\partial q_2 \partial q_1} & \dots & \dots \\ \dots & \dots & \dots \end{vmatrix} = |\text{Hess}| = \Delta q$$

Hesssche Determinante

$$\frac{D(p, \dots)}{D(q, \dots)} = \Delta q$$

$$e^{\frac{L^0}{2}} \dots$$

Notizen zum ...

$$\frac{1}{2} = \int \dots \frac{q-E}{2} \dots$$

$$= \frac{\int \dots e^{-\frac{q-E}{2\theta}} \dots}{\int \dots e^{-\frac{q-E}{2\theta}} \dots}$$

$$2 \text{ für } m = \dots$$

$$\frac{1}{2} \int \dots e^{-\frac{q^2}{2\theta}} \dots$$

$$\frac{1 \pm \sqrt{1+2\theta^3}}{\sqrt{1+2\theta}} = \frac{\theta}{2} = \frac{v_1}{2}$$

... ..

$$2 \sqrt{2\theta} \int \dots e^{-\frac{q-E}{2\theta}} \dots$$

Die Phasendichte als Integral:  $c \sim D(p, q_0)$  ...  
 $\int \dots e^{L^0} \dots$   
 $c \sim \int \dots e^{-\frac{q-E}{2\theta}} \dots$